## **Sentaurus™ Device QTX User Guide**

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# **SYNOPSYS®**

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## **Contents**





#### **[Chapter 2 Getting Started Example 17](#page-22-0)** 17





<span id="page-4-0"></span>This user guide describes the Synopsys® Sentaurus™ Device QTX tool that features the subband-based Boltzmann transport equation solver. This user guide is intended for users of TCAD Sentaurus involved in advanced CMOS transport modeling.

### <span id="page-4-1"></span>**Conventions**



The following conventions are used in Synopsys documentation.

### <span id="page-4-2"></span>**Customer Support**

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- support-tcad-jp@synopsys.com from Japan.

## <span id="page-6-0"></span>CHAPTER 1 Introduction to Sentaurus Device QTX

*This chapter describes the features of the subband-based Boltzmann transport equation solver of Sentaurus Device QTX and the physical models it uses for modeling quasi-ballistic transport.*

### <span id="page-6-1"></span>**Features of Sentaurus Device QTX**

Sentaurus Device QTX features the subband-based Boltzmann transport equation solver (hereafter, referred to as the *Subband-BTE solver*) that takes into consideration the effects of strong quantum confinement as well as quasi-ballistic transport along the channel of small FinFETs and nanowires. Both purely ballistic transport and scattering due to nonpolar phonon, polar optical phonon, surface roughness, alloy, and Coulomb scattering can be treated.

The subbands used by the Subband-BTE solver are provided by the solution of the Schrödinger equation on 2D slices along the channel. The solution of the Subband-BTE solver provides the distribution functions of carriers in each subband. This information, combined with the subband dispersion and the wavefunctions from the Schrödinger solver, provides the carrier density used in the Poisson equation. As shown in [Figure 1](#page-6-2), at each bias point, the Schrödinger equation, the subband-based Boltzmann transport equation (subband-BTE), and the Poisson equation are iterated until convergence.



<span id="page-6-2"></span>Figure 1 Schematic of the iteration procedure for the self-consistent solution of the Schrödinger equation, the subband-BTE, and the Poisson equation

The main output of the solution of the subband-BTE is the current flowing through the channel. In addition, several internal quantities related to the solution of the subband-BTE, such as the distribution function, the current spectrum, and the carrier velocity, can be saved to TDR files.

### <span id="page-7-0"></span>**Similarities to Sentaurus Band Structure**

The Subband-BTE solver shares many commands and physical models with Sentaurus Band Structure.

For descriptions of common commands, models, and parameters, refer to Part II of the *Sentaurus™ Device Monte Carlo User Guide*.

### <span id="page-7-1"></span>**Device Structure**

The simulation of carrier transport using the solution of the subband-BTE can be performed only for 3D devices. The use of Sentaurus Device QTX is most appropriate for devices with a well-defined channel, in which geometric quantum confinement plays a strong role. This is typically the case in small FinFETs and nanowire devices. The device structure is loaded from a TDR file using the LoadDevice command.



<span id="page-7-2"></span>Figure 2 Three-dimensional nanowire device showing slices and contacts

### <span id="page-8-0"></span>**Requirements and Restrictions on Device Structure and Mesh**

Sentaurus Device QTX has some requirements and restrictions on the type of device structure and mesh that can be treated, namely:

- The 3D mesh must be composed of all tetrahedra.
- The mesh within the channel should essentially be an extruded mesh from a 2D cross section. That is, along the channel direction, the mesh points should all fall on the slice locations.

### <span id="page-8-1"></span>**Device Axes and Orientation**

The device axis along the channel direction must be the z-axis. It must be set up correctly during structure generation. The confinement axes must be the x-axis and the y-axis. The orientation of the device relative to the crystallographic axes can be specified using the xDirection and yDirection parameters of the Physics command.

### <span id="page-8-2"></span>**Contacts**

It is expected that the left and right ends of the channel are terminated by contacts, as shown in [Figure 2 on page 2](#page-7-2). Additional contacts, typically one gate contact, can be included in the structure as well. However, the current is computed only for the left and right contacts of a channel.

During the solution of the Poisson equation, a special Neumann boundary condition should be used at the left and right contacts to enable charge neutrality there to be maintained, even in the presence of quantum confinement. To enable this boundary condition, the Physics contact command can be used along with type=Neumann, for example:

```
Physics contact=source type=Neumann
Physics contact=drain type=Neumann
```
A lumped resistance in units of ohm  $(\Omega)$  can be specified for the left and right contacts of a sliced channel. For example, to specify a lumped resistance of  $1 k\Omega$  for contacts named source and drain, specify:

```
Physics contact=source resist=1.0e3
Physics contact=drain resist=1.0e3
```
### <span id="page-9-0"></span>**Sliced Channels**

To solve the subband-BTE along the channel of a device, the subbands from the solution of the 2D Schrödinger equation at slices along the channel must be computed first. A group of these slices forms a so-called sliced channel. On each slice, a 2D Schrödinger equation is solved for the subband dispersions and wavefunctions. Based on these quantities, along with the distribution function for each subband, the carrier density on each node in the slice can be computed. Then, this is interpolated on to the 3D device mesh for solving the 3D Poisson equation.

### <span id="page-9-1"></span>**Creating a Sliced Channel**

You create a sliced channel using the [Math](#page-30-1) slicedChannel command (see Math [slicedChannel on page 25\)](#page-30-1). This is similar to the Math nonlocal command used in other parts of Sentaurus Band Structure. The parts of the device structure to be included in the sliced channel can be specified by using either a list of regions, or a bounding box, or both.

For example, to create a sliced channel consisting of only the silicon region, named sil, of a silicon nanowire, use the following command:

Math slicedChannel name=channel1 regions=[list si1]

During the creation of a sliced channel using the Math command, two key operations are performed automatically:

- First, the 2D slices along the channel are extracted automatically.
- Second, a nonlocal area is created on each slice with the name given by the sliced channel name appended by .NL, for example, channel1.NL.

### <span id="page-9-2"></span>**Specifying a Schrödinger Solver for a Sliced Channel**

Specifying a Schrödinger solver for a sliced channel is very similar to specifying a Schrödinger solver for a nonlocal area for mobility calculations in Sentaurus Band Structure. For a sliced channel, the Physics slicedChannel command is used, and the specified Schrödinger solver is used for all slices in the sliced channel (see [Physics slicedChannel on page 27\)](#page-32-1).

Examples:

```
Physics slicedChannel=channel1 eSchrodinger=Parabolic \
  valleys=[list Delta1 Delta2 Delta3] Nsubbands=8 Nk=16 Kmax=0.6 \
  a0=5.43e-8 correction=3
```

```
Physics slicedChannel=channel1 hSchrodinger=6kp \
  valleys=[list Gamma] Nsubbands=32 Nk=16 Kmax=0.6 a0=5.43e-8
```
### <span id="page-10-0"></span>**Model and Parameter Specification**

Similar to the specification of models and parameters for the subband and mobility calculations in Sentaurus Band Structure, for the Subband-BTE solver, models and parameters can be specified on a material or region basis.

These specified models and parameters will be used on all slices that contain the specified material or region. In the following sections, two model types are highlighted that are particularly important for the Subband-BTE solver: valley models and scattering models.

### <span id="page-10-1"></span>**Valley Models**

Valley models are used to specify the band model and parameters to use when solving the Schrödinger equation.

The following valley models can be used for the subband-BTE: ConstantEllipsoid, 2kpEllipsoidal, 6kp, and 3kp. The first two models can be used for the description of electrons, the 6kp valley model gives an accurate description of holes, and the 3kp valley model is an approximation of 6kp, where the split-off band is not captured due to the neglect of spin-orbit coupling.

### <span id="page-10-2"></span>**Scattering Models**

Scattering models are used to model specific, microscopic scattering transitions such as phonon and surface roughness scattering. All of the scattering models that can be used for mobility calculations in Sentaurus Band Structure on 2D structures can also be used when solving the subband-BTE. These models include elastic and inelastic nonpolar phonon scattering, polar optical phonon scattering, alloy scattering, surface roughness scattering, and Coulomb scattering from bulk impurities, interface charge, and traps. In addition, the surface roughness and the Coulomb scattering models can be screened using either the scalar or tensor dielectric function models.

A standard set of scattering models is set up for silicon regions by default, as described in the documentation of Sentaurus Band Structure in the *Sentaurus™ Device Monte Carlo User Guide*.

For purely ballistic simulations, all scattering models must be removed. To remove all scattering models, use the following command:

Physics material=all ScatteringModel removeAll

#### <span id="page-11-0"></span>RTA Scattering Model

The relaxation time approximation (RTA) scattering model is available for solving the subband-BTE. This model is implemented as an intra-subband scattering mechanism using the transition rate given in  $[1]$ :

$$
S_{\text{RTA}} = \frac{f_n(z, \varepsilon) - f_n^{\text{leq}}(z, \varepsilon)}{\tau}
$$
 (1)

where  $\varepsilon$  is the total energy and  $f_n^{\text{eq}}$  is an equilibrium Fermi–Dirac distribution function defined in terms of a local Fermi level,  $E_F^{\text{eq}}$ : ε is the total energy and  $f_n^$ <sup>n</sup> leq  $E_F^{\text{leq}}$ 

$$
f_n^{\text{leq}}(z,\varepsilon) = \frac{1}{1 + e^{(\varepsilon - E_F^{\text{leq}})/k_B T}}
$$
(2)

where  $T$  is the ambient temperature. The only adjustable parameter in the RTA scattering model is  $\tau$ , the relaxation time which is taken as a constant. The local Fermi level,  $E_F^{\text{leq}}$ , is determined by enforcing a carrier conservation condition on the equilibrium Fermi–Dirac distribution function integrated over energy. This produces an additional equation that must be solved along with the subband-BTE. The local Fermi level can be saved to a TDR file for visualization using the BTESaveCC command (see [BTESaveCC on page 34\)](#page-39-1).

The RTA scattering model serves two purposes:

- First, the RTA scattering model helps to stabilize the numeric solution of the subband-BTE. By default, an RTA scattering model is always used with a built-in  $\tau$  = 1e-10 s, which is a long enough relaxation time so as not to impact the computed current and is a short enough relaxation time so as to provide efficient numeric stability. This value of  $\tau$  for the RTA scattering model can be changed when selecting the numeric Subband-BTE solver as described next (see [Eq. 3](#page-12-2)).
- Second, the RTA scattering model provides an efficient model that can be used to mimic more complicated mechanisms that might incur significant runtime or are not yet available. For this purpose, additional RTA scattering models can be specified using the Physics ScatteringModel command on a per-region and per-valley basis, for example:

```
Physics region=si1 ScatteringModel=RTA name=rta1 valleys=[list Delta1] \
  tau=1.0e-13
```
<span id="page-12-2"></span>The final value of  $\tau$  that is used in the subband-BTE is computed using Mathiessen's rule considering the built-in RTA scattering model and all user-defined RTA scattering models for a particular region and valley, that is:

$$
\frac{1}{\tau} = \frac{1}{\tau_{\text{built-in}}} + \sum_{i} \frac{1}{\tau_i} \tag{3}
$$

#### <span id="page-12-0"></span>Surface Roughness

Surface roughness scattering is activated by specifying the SRFor2D scattering model. This model is used because it is applied to 2D slices. The SRFor2D model applies surface roughness scattering only at semiconductor–insulator interfaces.

During the creation of the sliced channel, Sentaurus Device QTX automatically extracts any insulator regions that are adjacent to the specified semiconductor regions. This does not affect the solution of the Schrödinger equation on the sliced channel, but enables the SRfor2D model to work correctly even if no insulator regions are specified explicitly.

### <span id="page-12-1"></span>Coulomb Scattering

When used with Sentaurus Device QTX, Coulomb scattering can be applied to both bulk doping and interface charge.

To speed up the calculation of the matrix elements for Coulomb scattering, the whenToCompute parameter can be used to specify when the matrix elements should be recomputed. The options for this parameter are:

- The EveryIteration option causes the matrix elements to be evaluated for each Poisson Newton iteration. This option provides the best accuracy but requires the most runtime. This is the default value.
- The FirstIteration option causes the matrix elements to be evaluated only for the first Poisson Newton iteration during a Solve command. Essentially, the solution from the previous bias point is used to compute the matrix elements. For subsequent Newton iterations, the matrix elements are reused when computing the scattering rate. This approach can potentially affect accuracy, but it greatly reduces the runtime. For bias ramps, this approach has worked well. If accuracy is important, a bias point can be repeated.

The following example shows how to specify the whenToCompute parameter:

```
Physics material=Silicon ScatteringModel=Coulomb name=eCoulomb \
  transitionType=Intravalley valleys=[list Delta1 Delta2 Delta3] \
  screening=Lindhard whenToCompute=FirstIteration
```
#### <span id="page-13-0"></span>Screening

Some of the scattering models, such as surface roughness and Coulomb, can be screened using one of the available screening models. By default, no screening is used. The required screening model can be selected using the screening parameter for each scattering model.

This parameter can be set to one of the following values:

- none for no screening
- Lindhard for the scalar Lindhard screening model
- LindhardTDF for the tensor Lindhard screening model

The screening models utilize the polarization  $(\Pi)$ , which characterizes the strength of the screening, and the dielectric form-factor  $(F)$ , which characterizes the screening effectiveness of each subband pair in terms of the wavefunction overlap with the Coulomb Green's function.

The polarization is given by:

$$
\Pi_{nn'}(q) = \frac{g_s \cdot g_v}{2\pi} \int dk \frac{f_{n'}(k+q) - f_n(k)}{E_n(k) - E_{n'}(k+q)}
$$
(4)

where:

- $\bullet$   $g_s$  and  $g_v$  are the spin and valley degeneracy, respectively.
- $E_n$  and  $f_n$  are the dispersion and distribution function for subband *n*, respectively.

The dielectric form-factor is given by:

$$
F_{mm'}^{nn'}(q) = \iint dx dx' \ \Psi_m(x) \cdot \Psi_{m'}^{\dagger}(x) G(q, x, x') \Psi_n^{\dagger}(x') \cdot \Psi_{n'}(x')
$$
 (5)

where  $\Psi_m$  is the wavefunction for subband m, and G is the Coulomb Green's function.

#### <span id="page-13-1"></span>Tensor Dielectric Function

The tensor dielectric function (TDF) in the TDF screening model is given by:

$$
\varepsilon_{mm'}^{nn'}(q) = \delta_{mn}\delta_{m'n'} + \Pi_{nn'}(q)F_{mm'}^{nn'}(q)
$$
\n(6)

The tensor Lindhard screening model (LindhardTDF) is time consuming to compute, particularly when a large number of subband pairs is used for the screening calculation. The number of subband pairs can be controlled using the tdfDegenTol parameter of the Physics slicedChannel command when specifying the Subband-BTE solver. Subband pairs are used in the screening calculation if their subband edges are within tdfDegenTol of each other. It is recommended to set this to a value of 1.0e-4 eV.

#### <span id="page-14-0"></span>Scalar Dielectric Function

The scalar dielectric function (SDF) in the SDF screening model is given by:

$$
\varepsilon(q) = 1 + \sum_{n} \Pi_{nn}(q) F_{nn}^{nn}(q)
$$
\n(7)

The sum is over the subbands. While it is expensive to compute, the SDF screening model is much faster than the TDF screening model and provides a good trade-off between accuracy and runtime.

#### <span id="page-14-1"></span>Runtime Speedup With Screening

The screening models are computationally expensive. They are heavily parallelized and you should use as many threads as possible. However, even with many threads, for a structure with many slices, the overall wallclock time when using screening may still be quite long.

To provide a trade-off between accuracy and runtime, when specifying the Subband-BTE solver, the whenToComputeScreening parameter of the Physics slicedChannel command can be used to control how often the screening model is evaluated. The options for this parameter are:

- The EveryIteration option causes the screening model to be evaluated for each Poisson Newton iteration. This option provides the best accuracy but requires the most runtime. This is the default value.
- The FirstIteration option causes the screening model to be evaluated only for the first Poisson Newton iteration during a Solve command. Essentially, the solution from the previous bias point is used to compute the screening dielectric function. For subsequent Newton iterations, the dielectric function is reused. This approach can potentially affect accuracy, but it greatly reduces the runtime. For bias ramps, this approach has worked well. If accuracy is important, a bias point can be repeated.

The following example shows how to specify the whenToComputeScreening parameter and the tdfDegenTol parameter:

```
Physics slicedChannel=channel1 eBTESolver=Numeric tdfDegenTol=1.0e-4 \
  whenToComputeScreening=FirstIteration
```
#### <span id="page-14-2"></span>Limiting the Subband Scattering Range

For transition types such as Intravalley, Intervalley, and gIntervalley, a state in the initial subband is allowed, by default, to scatter to all valid final states in all valid final subbands. When using a large number of subbands and a fine energy grid for the solution of the subband-BTE, this produces a very large system matrix that can be very time consuming to solve. Restricting the number of final subbands to which an initial state scatters can help to reduce the overall system size and to improve runtime performance. This can be achieved by using the subbandScatteringRange parameter of the Physics slicedChannel command (see [Physics slicedChannel on page 27](#page-32-1)).

### <span id="page-15-0"></span>**Subband-Based Boltzmann Transport Equation**

With quantum confinement in two dimensions, the subband-BTE is reduced to one dimension along the channel within each subband. The solution of the subband-BTE gives the distribution function as a function of the  $k$ -vector and the channel coordinate for each subband. Using  $z$ for the channel axis, for subband  $n$ , the subband-BTE is  $[2]$ :

$$
-\frac{\partial f_n}{\partial k} \frac{1}{\hbar} \frac{\partial E_n}{\partial z} + \frac{\partial f_n}{\partial z} \frac{1}{\hbar} \frac{\partial E_n}{\partial k} = S_{n, \text{ in}} - S_{n, \text{ out}}
$$
(8)

where:

- **f**<sub>n</sub> is the distribution function for subband *n*.
- $E_n$  is the dispersion for subband *n*.

With Fermi–Dirac statistics, the in-scattering term  $(S_{n, in})$  and the out-scattering term  $(S_{n, out})$ are given by:

$$
S_{n, \text{ in}} = \sum_{n'} \frac{1}{2\pi} \int dk' S_{n'n}(k', k) f_{n'}(z, k') [1 - f_n(z, k)]
$$
  
\n
$$
S_{n, \text{ out}} = \sum_{n'} \frac{1}{2\pi} \int dk' S_{nn'}(k, k') f_n(z, k) [1 - f_n'(z, k')]
$$
\n(9)

where  $S_{n'n}(k', k)$  is the total transition rate due to scattering.

As boundary conditions on the distribution functions, equilibrium Fermi–Dirac distribution functions at the left and right contacts are applied to carriers being injected into the channel.

### <span id="page-15-1"></span>**Selecting a Subband-BTE Solver**

Two Subband-BTE solvers are available for solving the subband-BTE.

The faster and simpler solver is called AnalyticBallistic. As its name suggests, it can be used only for purely ballistic simulations. This solver uses the analytic solution to the subband-BTE, which is determined by projecting the injecting, equilibrium Fermi–Dirac distribution functions from the contacts throughout the device as determined by the top of each subband barrier.

The second solver is called Numeric. It performs an actual numeric solution to the subband-BTE. This solver can be used for purely ballistic transport as well as with all of the available scattering models including the RTA scattering model.

For example, to activate the AnalyticBallistic subband-BTE for electron transport for a sliced channel named channel1, you can use the following command:

Physics slicedChannel=channel1 eBTESolver=AnalyticBallistic

For a subband-BTE simulation of holes, you can use the hBTESolver argument:

```
Physics slicedChannel=channel1 hBTESolver=AnalyticBallistic
```
### <span id="page-16-0"></span>**Energy Grid**

The subband-BTE is solved on a 2D tensor-product grid. One axis of this grid is along the channel direction and is referred to as the *channel coordinate axis*. The grid points along this axis are given by the slice locations along the channel. The other axis of the BTE solution grid is the *total energy axis*. Here, *total energy* refers to the total energy of a carrier in a particular subband as determined by the sum of the minimum subband energy and the kinetic energy of the carrier.

The minimum and maximum energies used in the energy grid are determined automatically. The nominal energy grid spacing is set using the deltaE argument of the Physics slicedChannel command when selecting a Subband-BTE solver (see [Physics slicedChannel](#page-32-1) [on page 27\)](#page-32-1).

In addition, a special refinement around the top of each subband can be activated to improve the accuracy of the calculation and to improve convergence. This is activated using the resolveTOBEnergy argument. The smallest energy spacing used for this refinement is set using the minDeltaE argument.

For example, a typical energy grid setup for a sliced channel named channel 1 is:

```
Physics slicedChannel=channel1 eBTESolver=Numeric deltaE=4e-3 \
  resolveTOBEnergy=1 minDeltaE=1.0e-4
```
### <span id="page-17-0"></span>**Numeric Options**

When specifying the Subband-BTE solver (see [Physics slicedChannel on page 27\)](#page-32-1), several options related to numeric details can be set, in particular:

■ useParabolicFitForToB=*Boolean*

When true, this argument activates a special algorithm to more accurately locate the top of the source–drain barrier. This mainly applies to FET devices. Default: true.

■ approximateJacobian=*Boolean*

When true, this argument activates a special algorithm to approximate the Jacobian that is used when solving the BTE system. This can be used to greatly reduce memory consumption with a small impact on convergence. Default: true.

■ useKdependentWFForBTEDensity=*Boolean*

When true, this argument computes the carrier density based on  $k$ -dependent wavefunctions. When this argument is false,  $\Gamma$ -point wavefunctions are used. By default, this argument is deactivated to improve convergence.

- **NOTE** If useKdependentWFForBTEDensity=0 and useKdependentWF=1 are set when specifying the Schrödinger solver, *k*-dependent wavefunctions will still be used when computing the scattering rates.
- **NOTE** If reorderDispersion is switched on, you should also switch on useKdependentWFForBTEDensity and useKdependentWF to have good convergence behavior.
- NkForScreening=*Integer*

This argument sets the number of  $k$ -points to use when evaluating screening. Default: 32.

### <span id="page-17-1"></span>**Solving at One Bias**

The solution of the subband-BTE at a specific bias is initiated using the Solve command. Similar to calculations using Sentaurus Band Structure, the bias on each contact is set using the V(<contact>) syntax, where <contact> refers to the contact name. For example, to initiate a solve at specific biases on the source, drain, and gate contacts, use the following command:

```
Solve V(gate) = 1.0 V(drain) = 1.0 V(source) = 0.0
```
If the bias on a particular contact is not specified explicitly in the Solve command, the previously specified bias is used. By default, the bias on all contacts is set to 0 at the start of the simulation.

### <span id="page-18-0"></span>**Convergence Criteria**

During the solution of the subband-BTE, three equations are solved in an iterative fashion: the Schrödinger equation, the subband-BTE, and the Poisson equation. The overall iteration of these equations is divided into an outer iteration of the Poisson equation and an inner iteration of the subband-BTE after a single solve of the Schrödinger equation. This is shown in [Figure 1](#page-6-2) [on page 1](#page-6-2).

The convergence criteria for the outer Poisson equation are set by the currentConvTol and potentialUpdateTolerance arguments of the Math command. The argument potentialUpdateTolerance, set in units of  $k_B T/q$ , determines the maximum-allowed change in the potential below which the Poisson equation is considered to have converged. For the subband-BTE system of equations, it is recommended to set this argument to a value of approximately 3.0e-3.

The currentConvTol argument determines the relative change in terminal currents between subsequent Poisson iterations below which the subband-BTE system is considered to have converged.

The subband-BTE system is considered to have converged when either of these two convergence criteria are met. For example, if potentialUpdateTolerance=3.0e-3 and currentConvTol=1.0e-3, a bias point for the subband-BTE will be considered converged if the maximum change in the potential at a particular Poisson Newton iteration is below  $3.0e-3 \times (k_B T/q)$  or the change in terminal current relative to the previous Newton iteration is below 1.0e-3. By default, currentConvTol is set to 0.0, meaning that convergence is completely controlled by the potentialUpdateTolerance argument.

The convergence criteria of the inner solution of the subband-BTE are set using the following arguments:

- distFuncUpdateTolerance for the distribution functions.
- rtaUpdateTolerance for the local Fermi level used in the RTA scattering model.

It is recommended to use the default values for these arguments.

The solution of the subband-BTE requires an initial guess for the distribution functions. For the first several Poisson iterations, this guess is supplied by solving the subband-BTE using only the RTA scattering model. After these initial Poisson iterations, the initial guess is supplied by the solution of the subband-BTE from the previous Poisson iteration. The exact number of Poisson iterations for which the RTA solution is used as an initial guess is set by the iterationsForRTAGuess parameter of the Math command. It is recommended to use its default value of 5.

For details about these arguments, see [Math on page 23](#page-28-3).

### <span id="page-19-0"></span>**Output**

The primary output of the solution of Sentaurus Device QTX is the current, in ampere, flowing through the left and right contacts at the end of the channel. These current values are printed to the screen at the end of a Solve command, and the values also are stored in the bias log file. Furthermore, in the bias log file, when a nonzero lumped resistance is specified, in addition to the applied bias that is stored under  $V(\text{contact}>)$ , the so-called inner voltage on the inner end of the lumped resistor is stored under Vinner (<contact>).

In addition to the current, several different TDR files with data from the solution of the Subband-BTE solver can be saved. These are described in the next sections. See [Chapter 3 on](#page-28-4) [page 23](#page-28-4) for full details. [Chapter 2 on page 17](#page-22-2) presents an example.

### <span id="page-19-1"></span>**2D TDR File With Fields Over Channel Coordinate–Energy Space**

You can use the BTESaveE command to save a 2D TDR file with fields that are defined on the channel coordinate–energy grid. These fields include quantities such as the distribution function, the current spectrum, and the density-of-states (DOS) in each subband, or in each valley, or in total (see [BTESaveE on page 30\)](#page-35-1).

### <span id="page-19-2"></span>**2D TDR File With Fields Over Channel Coordinate– k-Space**

You can use the BTESaveK command to save a 2D TDR file with fields that are defined over channel coordinate– $k$ -space (see [BTESaveK on page 32\)](#page-37-1).

### <span id="page-19-3"></span>**TDR (XY) File With Fields Versus Channel Coordinate**

You can use the BTESaveCC command to save a TDR (xy) file with fields versus the channel coordinate. The fields that can be saved include the total inversion charge, the total current, the total carrier velocity, as well as these quantities per subband. In addition, the occupancy of each subband can be saved (see [BTESaveCC on page 34\)](#page-39-1).

### <span id="page-20-0"></span>**Slice-Related Fields**

Each slice that is used in the solution of the subband-BTE is basically a 2D cross section. Similar to the direct treatment of a 2D device structure using Sentaurus Band Structure, various fields across this 2D cross section, as well as fields in  $1D$   $k$ -space, can be saved using two commands.

### <span id="page-20-1"></span>2D TDR File With Fields Over XY Real Space

You can use the SaveSlice command to save quantities from the Schrödinger solver on a slice such as the subband energy and the wavefunctions, as well as quantities such as the conduction band energy (see [SaveSlice on page 36](#page-41-1)).

### <span id="page-20-2"></span>TDR (XY) File With Fields Over 1D k-Space

You can use the  $SavesliceK$  command to save the  $1D$   $k$ -space dispersion for each subband computed by the Schrödinger solver (see [SaveSliceK on page 37](#page-42-1)).

### <span id="page-20-3"></span>**References**

- <span id="page-20-4"></span>[1] S. Jin, T.-W. Tang, and M. V. Fischetti, "Simulation of Silicon Nanowire Transistors Using Boltzmann Transport Equation Under Relaxation Time Approximation," *IEEE Transactions on Electron Devices*, vol. 55, no. 3, pp. 727–736, 2008.
- <span id="page-20-5"></span>[2] D. Esseni, P. Palestri, and L. Selmi, *Nanoscale MOS Transistors: Semi-Classical Transport and Applications*, Cambridge: Cambridge University Press, 2011.

**1: Introduction to Sentaurus Device QTX** References

<span id="page-22-2"></span><span id="page-22-0"></span>*This chapter presents an example that uses Sentaurus Device QTX to compute an*  $I_d$ *–V<sub>g</sub> curve for a silicon NMOS nanowire.* 

This chapter describes the command file for simulating the ballistic transport in a silicon nanowire device using Sentaurus Device QTX. The various sections of the command file used to load the device structure, to set up the models, to specify the sliced channel, and to perform the  $I_d-V_g$  simulation are described.

### <span id="page-22-1"></span>**Device Structure**

[Figure 3](#page-22-3) shows the device structure for this example. The device is a cylindrical silicon NMOS nanowire with a diameter of 5 nm, a gate length of 13 nm, and source and drain extension regions of length 10 nm. The source and drain are uniformly doped n-type to  $2e20 \text{ cm}^{-3}$ . This example computes an  $I_d-V_g$  curve at  $V_d = 0.6$  V in the ballistic limit using the Numeric Subband-BTE solver.



<span id="page-22-3"></span>

### <span id="page-23-0"></span>**Loading the Device Structure**

LoadDevice tdrFile=nanowire3D\_0.tdr

The device structure, nanowire3D  $0.tdr$ , is loaded using the LoadDevice command. Upon loading, the tool automatically sets up the default models and parameters in the silicon and oxide regions.

### <span id="page-23-1"></span>**Setting Up the Model and the Classical Solve**

```
# Simple model for hole density
Physics material=Silicon hBulkDensity=hFermiDensity Nv=3.10e19
# Remove all scattering models for ballistic simulation
Physics material=all ScatteringModel removeAll
# Use Neumann boundary condition on source and drain
Physics contact=source type=Neumann
Physics contact=drain type=Neumann
# Set workfunction
Physics contact=gate workfunction=4.25
# Set orientation for <110> channel
Physics xDirection=[list 1 1 0] yDirection=[list 0 0 1]
# Classical solve
Solve
```
These commands set up particular models and parameters, starting with specifying a simple model for the hole density. Then, all of the scattering models are removed so that a ballistic simulation can be performed. The boundary condition type for the source and drain contacts is set to Neumann. Finally, the gate workfunction is set to 4.25 V.

After the required models are specified, a classical solve of only the Poisson equation is performed in equilibrium, that is, zero bias on all contacts. A *classical solve* means that the Schrödinger equation is not used. The resulting classical solution will serve as the initial guess for the first solve using the Schrödinger equation.

### <span id="page-24-0"></span>**Specifying the Sliced Channel**

# Create sliced channel over all device for si1 region Math slicedChannel name=channel1 regions=[list si1] # Set up parabolic Schrodinger on the sliced channel Physics slicedChannel=channel1 eSchrodinger=Parabolic \ valleys=[list Delta1 Delta2 Delta3] Nsubbands=8 Nk=16 Kmax=0.6 \ a0=5.43e-8 correction=3 # Select Numeric Subband-BTE solver and set energy grid parameters Physics slicedChannel=channel1 eBTESolver=Numeric deltaE=4.0e-3 \ resolveTOBEnergy=1 minDeltaE=1.0e-4

This section of the command file specifies the sliced channel, the Schrödinger solver, and the Subband-BTE solver.

The Math slicedChannel command creates a sliced channel named channel1 over the region named si1 throughout the entire device.

The first Physics slicedChannel command specifies that a parabolic Schrödinger solver must be used on all slices of the sliced channel. This command refers to the Delta1, Delta2, and Delta3 valley models that were created by default during the LoadDevice command.

The second Physics slicedChannel command changes the Subband-BTE solver to Numeric and specifies some energy grid parameters.

### <span id="page-24-1"></span>**Performing the I<sub>d</sub>-V<sub>g</sub> Simulation**

```
# Set some convergence parameters. It continues to next bias if not converged
Math potentialUpdateTolerance=6.0e-3 iterations=30 doOnFailure=0
# Solve at equilibrium and then at required drain bias
Solve V(gate)=0.0 V(drain)=0.0Solve V(drain)=0.6
# Ramp gate using a Tcl foreach loop
foreach Vg [list 0.0 0.1 0.2 0.3 0.4 0.5 0.6] {
  Solve V(gate)=$Vg
  AddToLogFile name=Id value=[GetLast name=I(drain)]
}
```
This section of the command file computes the  $I_d-V_g$  curve. First, the Math command specifies the convergence tolerance for the Poisson equation. The number of allowed iterations for each bias is set to 30, and the doOnFailure argument is set such that the tool will continue to the next bias, even if the convergence criteria are not met.

This series of Solve commands will write data to the default bias log file named subbandBTE\_Quickstart.plt. An extra user-defined quantity named Id is added to the bias log file to make plotting the  $I_d-V_g$  curve easier in Sentaurus Visual. [Figure 4](#page-25-1) shows the resulting  $I_d-V_g$  curve.



<span id="page-25-1"></span>Figure 4 Ballistic  $I_d-V_g$  curve from the example

### <span id="page-25-0"></span>**Saving TDR Files**

```
# Save 3D TDR
Save tdrFile=subbandBTE Quickstart.tdr \
  models=[list DopingConcentration eDensity hDensity ConductionBandEnergy \
          eQuasiFermiEnergy]
# Save fields over ChannelCoord-Energy
BTESaveE tdrFile=subbandBTE Quickstart CC E.tdr \
  models=[list Delta1_0_DistributionFunction Delta1_0_CurrentSpectrum \
          Delta1_0_SubbandEnergy Delta1_0_DOS]
# Save distribution function over ChannelCoord-k
BTESaveK tdrFile=subbandBTE Quickstart CC K.tdr \
  models=[list Delta1_0_DistributionFunction]
```

```
# Save subband quantities over ChannelCoord
BTESaveCC tdrFile=subbandBTE Quickstart CC.tdr \
  models=[list NinvTotal CurrentTotal VelocityTotal \
          Delta1_0_Occupancy Delta3_0_Occupancy]
# Save fields over a slice
SaveSlice tdrFile=subbandBTE Quickstart Slice.tdr channelCoord=20.0e-3 \
  models=[list eDensity Delta1_0_Wavefunction]
# Save dispersion at a slice
SaveSliceK tdrFile=subbandBTE Quickstart SliceK.tdr channelCoord=20.0e-3 \
  models=[list Delta1_0_Dispersion]
```
This section of the command file writes various TDR files with different types of data from the solution of the subband-BTE:

- The Save command writes a 3D TDR file with the 3D device structure and the selected models.
- The BTESaveE command writes a 2D TDR file over channel coordinate–energy space. In this example. the distribution function, the current spectrum, the subband energy, and DOS for the Delta1 0 subband are included. [Figure 5 on page 22](#page-27-0) shows the plot of the resulting distribution function.
- **The BTESaveK command saves a TDR file over channel coordinate–** $k$ **-space with only the** distribution function for the Delta1\_0 subband.
- The BTESaveCC command saves a TDR (xy) file with the selected models versus the channel coordinate. [Figure 6 on page 22](#page-27-1) shows the plot of the resulting total inversion change (NinvTotal) and the carrier velocity (VelocityTotal) along the channel.
- **The final two commands, SaveSlice and SaveSliceK, save real-space and**  $k$ **-space** quantities, respectively, from the slice located closest to the channel coordinate of 20 nm.



<span id="page-27-0"></span>



<span id="page-27-1"></span>Figure 6 Plot of the total inversion charge and carrier velocity along the nanowire channel from the BTESaveCC command

## <span id="page-28-4"></span><span id="page-28-0"></span>CHAPTER 3 Commands of Sentaurus Device QTX

*This chapter summarizes the commands used by Sentaurus Device QTX.*

### <span id="page-28-1"></span>**Math and Physics Commands**

This section describes the commands used to specify the geometry and the physics of a sliced channel as well as the convergence criteria for the solution of the subband-BTE.

### <span id="page-28-3"></span><span id="page-28-2"></span>**Math**

Sets the convergence criteria for solving the subband-BTE.

#### **Syntax**

```
Math [currentConvTol=Double] [distFuncUpdateTolerance=Double] \
  [iterationsForRTAGuess=Integer] [potentialUpdateTolerance=Double] \
   [rtaUpdateTolerance=Double]
```
#### **Arguments**



#### **Description**

The basic Math command sets the convergence criteria for solving the subband-BTE using the Numeric Subband-BTE solver along with the RTA scattering model. The subband-BTE is considered converged when both criteria specified by distFuncUpdateTolerance and rtaUpdateTolerance are met.

The iterationsForRTAGuess parameter can be used to control the initial guess for the solution of the subband-BTE. Setting this to a large value can help to improve convergence at the beginning of a Solve; however, it can result in a large overall number of Poisson iterations.

#### **Examples**

```
Math distFuncUpdateTolerance=1.0e-7 rtaUpdateTolerance=1.0e-3
```
This command sets the update tolerance on the distribution functions to 1.0e-7 and on the RTA Fermi levels to 1.0e-3 eV.

### <span id="page-30-1"></span><span id="page-30-0"></span>**Math slicedChannel**

Creates a sliced channel within a 3D device.

#### **Syntax**

```
Math slicedChannel name=String [regions=List] [channelAxis=String] \
   [minX=Double] [maxX=Double] [minY=Double] [maxY=Double] [minZ=Double] \
   [maxZ=Double]
```
#### **Arguments**



#### **Description**

The Math slicedChannel command creates a sliced channel within the 3D device that has been already loaded using the LoadDevice command. This command automatically creates the 2D slices along the channel as well as a nonlocal area for each slice. The name of the nonlocal area is set automatically to <slicedChannel>.NL. For example, if the sliced channel is named channel1, the nonlocal will be named channel.NL. As in the Math nonlocal command, only mesh elements contained within the specified geometry are included.

#### **Examples**

Math slicedChannel name=channel1 regions=[list si1 ox1]

This command created a new sliced channel named channel1. Slices along the channel axis, which defaults to z, are extracted automatically within the si1 and ox1 regions.

### <span id="page-31-0"></span>**Physics contact**

Sets the properties of the specified contact.

#### **Syntax**

Physics contact=*String* [resist=*Double*] [type=*String*]

#### **Arguments**



#### **Description**

The Physics contact command sets the properties of the specified contact. For the Subband-BTE solver, the type argument can be used to set the required Neumann boundary condition for the solution of the subband-BTE on the left and right contacts of each sliced channel. A lumped resistance can be specified for a contact through which current flows using the resist argument.

#### **Examples**

Physics contact=source type=Neumann Physics contact=drain type=Neumann

These commands set the boundary condition type on the source and drain contacts to Neumann.

### <span id="page-32-1"></span><span id="page-32-0"></span>**Physics slicedChannel**

Specifies the Schrödinger solver or the Subband-BTE solver to use on a sliced channel.

#### **Syntax**

To specify the Schrödinger solver to use on all slices in a sliced channel, use:

Physics slicedChannel=*String* eSchrodinger=*String* ...

To specify the Subband-BTE solver to use on a sliced channel, use:

```
Physics slicedChannel=String (eBTESolver=String | hBTESolver=String) \
  [approximateJacobian=Boolean] [deltaE=Double] [minDeltaE=Double] \
   [NkForScreening=Integer] [reorderDispersion=Boolean] \
   [resolveTOBEnergy=Boolean] [subbandScatteringRange=Integer] \
   [tau=Double] [tdfDegenTol=Double] \
   [useKdependentWF=Boolean] [useKdependentWFForBTEDensity=Boolean] \
   [useParabolicFitForToB=Boolean] [whenToComputeScreening=String]
```
#### **Arguments**



#### **3: Commands of Sentaurus Device QTX**

Math and Physics Commands



#### **Description**

The Physics slicedChannel command is used to specify the Schrödinger solver or the Subband-BTE solver to use on a sliced channel, depending on which parameters are used.

To specify the Schrödinger solver to use on all slices in a sliced channel, use the eSchrodinger or hSchrodinger argument. The use of this argument is identical to setting up a Schrödinger solver using the Physics nonlocal command, and all of the Schrödingerrelated arguments can be used here as well.

To specify the Subband-BTE solver to use on a sliced channel, use the eBTESolver or hBTESolver argument.

#### **Examples**

```
Physics slicedChannel=channel1 eSchrodinger=Parabolic \
  valleys=[list Delta1 Delta2 Delta3] Nsubbands=8 Nk=16 Kmax=0.6 a0=5.43e-8
```
This command specifies that a parabolic Schrödinger solver must be used on all slices of the sliced channel named channel1. The usual set of arguments for setting up a Schrödinger solver is specified as well.

Physics slicedChannel=channel1 eBTESolver=AnalyticBallistic deltaE=5.0e-3

This command specifies that the Subband-BTE solver for channel1 should change to the AnalyticBallistic Subband-BTE solver and that the nominal energy grid spacing should change to 5 meV.

### <span id="page-34-0"></span>**Saving TDR Files**

Several types of data and TDR files can be saved from Sentaurus Device QTX. Each type of TDR file is saved using a command specific to the type of data as described here.

### <span id="page-34-1"></span>**Specifying Subbands and Subband IDs**

Many of the quantities are defined on a subband basis. Therefore, you must specify exactly which subband and which quantity are required.

A particular subband is uniquely specified by its so-called subband ID, which is determined by the valley name of the subband and its subband index within this valley, separated by an underscore  $(\_)$ . For example, the subband ID of Deltal  $\circ$  corresponds to the 0th subband of the Delta1 valley.

**NOTE** Subband indexing starts from 0.

Therefore, a particular subband-based quantity is specified by giving the subband ID and then the quantity name, again separated by an underscore. For example, the name Delta1\_0\_DistributionFunction specifies the distribution function for the 0th subband of the Delta1 valley.

In the following command sections, various subband-based quantities are described using names such as <subbandID> DistributionFunction. Here, <subbandID> should be replaced by the subband ID such as Delta1 0.

### <span id="page-35-1"></span><span id="page-35-0"></span>**BTESaveE**

Saves the solution of the subband-BTE performed over the channel coordinate–energy space.

#### **Syntax**

BTESaveE tdrFile=*String* models=*List* [slicedChannel=*String*]

#### **Arguments**



The available quantities that can be included in the list of models to save are:





#### **Description**

The solution of the subband-BTE is performed over the channel coordinate–energy space. A few major quantities are solved or computed on this grid on a per-subband basis. The BTESaveE command allows you to save a 2D TDR file over the channel coordinate–energy space for these quantities.

For quantities other than SubbandEnergy, the subbandID can be used to save a quantity either per subband, or per valley, or as the total value summed over all subbands. To save a quantity per valley, the subbandID must contain only the valley name. To save the total value of a quantity, do not specify the subbandID. For example:

■ To save the DensitySpectrum for the Delta1\_0 subband, specify:

Delta1\_0\_DensitySpectrum

■ To save the sum of the DensitySpectrum for all subbands in the Delta1 valley, specify:

Delta1\_DensitySpectrum

To save the sum of the DensitySpectrum over all subbands, specify:

DensitySpectrum

#### **Examples**

```
BTESaveE tdrFile=FieldsOver CC E.tdr \
  models=[list Delta1_0_DistributionFunction Delta1_0_CurrentSpectrum]
```
This command saves a TDR file named FieldsOver\_CC\_E.tdr containing the distribution function and the current spectrum for the Delta1\_0 subband.

### <span id="page-37-1"></span><span id="page-37-0"></span>**BTESaveK**

Saves a 2D TDR file over the channel coordinate– $k$ -space for the specified quantities.

#### **Syntax**

BTESaveK tdrFile=*String* models=*List* [slicedChannel=*String*] [Nk=*Integer*] \ [Kmax=*Double*]

#### **Arguments**



The available quantities that can be included in the list of models to save are:



#### **Description**

When the solution of the subband-BTE is performed over the channel coordinate–energy space, it is sometimes helpful to visualize some quantities over the channel coordinate- $k$ space instead, in particular, the distribution function. The BTESaveK command saves a 2D TDR file over channel coordinate– $k$ -space for the specified quantities. The Nk and Kmax arguments are provided to specify the  $k$ -grid that should be used.

#### **Examples**

```
BTESaveK tdrFile=FieldsOver CC K.tdr \
  models=[list Delta1_0_DistributionFunction] Nk=101 Kmax=0.3
```
This command saves a TDR file named FieldsOver\_CC\_K.tdr containing the distribution function for the  $\text{Delta}_0$  subband over channel coordinate– $k$ -space. The  $k$ -grid that is saved extends from  $-0.3^* 2\pi/a_0$  to  $0.3^* 2\pi/a_0$  with 101 points.

### <span id="page-39-1"></span><span id="page-39-0"></span>**BTESaveCC**

Saves quantities from the solution of the subband-BTE that depend only on the channel coordinate.

#### **Syntax**

BTESaveCC tdrFile=*String* models=*List* [slicedChannel=*String*]

#### **Arguments**



The available quantities that can be included in the list of models to save are:





#### **Description**

Several quantities from the solution of the subband-BTE depend only on the channel coordinate. These include the inversion charge, the current, and the carrier velocity. These quantities can be saved on a per-subband basis. In addition to the quantities mentioned, the total value from all subbands can be saved.

#### **Examples**

BTESaveCC tdrFile=FieldsOver\_CC.tdr models=[list NinvTotal Delta1\_0\_Velocity]

This command saves a TDR (xy) file named FieldsOver\_CC.tdr containing the total inversion charge along the channel as well as the velocity in the 0th subband of the Delta1 valley. Since no sliced channel is specified, the first defined sliced channel is used automatically.

### <span id="page-40-0"></span>**Saving Slice-Related Fields**

Each 2D slice is like a 2D device with a 2D nonlocal area defined. Like a 2D device, real-space models over this 2D device can be saved as well as  $1D k$ -space data similar to what is done for Schrödinger and mobility calculations on a single 2D device in Sentaurus Band Structure.

The model syntax is exactly the same. The key difference here is that a particular 2D slice from a sliced channel must be selected. A 2D slice has a unique *address* given by its sliced channel name and its channel coordinate. The user-specified ChannelCoord snaps to the nearest actual slice position.

### <span id="page-41-1"></span><span id="page-41-0"></span>**SaveSlice**

Saves real-space models over a 2D slice to a TDR file.

#### **Syntax**

```
SaveSlice tdrFile=String channelCoord=Double models=List \
   [slicedChannel=String]
```
#### **Arguments**



#### **Description**

This command saves real-space models over the 2D slice specified with the slicedChannel and channelCoord arguments. The coordinates of the real-space mesh are saved in units of micrometer.

#### **Examples**

```
SaveSlice tdrFile=slice.tdr slicedChannel=channel1 channelCoord=20.0e-3 \
  models=[list ConductionBandEnergy Delta1_0_Wavefunction]
```
This command saves a 2D TDR file named slice.tdr containing the mesh and specified models over the 2D slice located in the sliced channel named channel1 and closest to the slice channel coordinate of  $20.0e-3 \mu m$ . Two models are saved: the relaxed conduction band energy and the norm of the wavefunction for the Delta1\_0 subband.

### <span id="page-42-1"></span><span id="page-42-0"></span>**SaveSliceK**

Saves a TDR (xy) file of  $k$ -space models over 1D  $k$ -space for the slice specified.

#### **Syntax**

```
SaveSliceK tdrFile=String channelCoord=Double models=List \
   [slicedChannel=String] [Nk=Integer] [Kmax=Double]
```
#### **Arguments**



#### **Description**

This command saves a TDR (xy) file of  $k$ -space models over 1D  $k$ -space for the slice specified with the slicedChannel and channelCoord arguments. The  $1D$   $k$ -space coordinates are saved in units of  $2\pi/a_0$ . The k-space grid that is used is specified using the Nk and Kmax arguments.

#### **Examples**

SaveSliceK tdrFile=slice K.tdr slicedChannel=channel1 channelCoord=20.0e-3 \ models=[list Delta1\_0\_Dispersion]

This command saves a TDR (xy) file named  $\text{slice\_K}. \text{tdr}$  containing a 1D  $k$ -space grid and specified models over 1D  $k$ -space for the slice located in the sliced channel named channel 1 and closest to the slice channel coordinate of  $20.0e-3 \mu m$ . In this example, only the dispersion for the Delta1\_0 subband is saved. The default values for Nk and Kmax are used.

**3: Commands of Sentaurus Device QTX** Saving Slice-Related Fields